

**Step 1:** Measured spectrogram  $I_0(p, \tau)$

$$I_0(p, \tau) = \left| \int_{-\infty}^{\infty} E_X(t - \tau) d(p + A(t)) e^{-i\varphi(p,t)} e^{\left(\frac{p^2}{2} + I_p\right)t} dt \right|^2 \quad (1)$$

$$\varphi(p, t) = \int_t^{\infty} \left[ pA(t') + \frac{A^2(t')}{2} \right] dt' \quad (2)$$

Experimentally,  $I_0(p, \tau)$  is an  $N_\epsilon \times N_\tau$  2-D array with  $N_\epsilon$  energy points and  $N_\tau$  delay points. Usually, the spectrogram is scanned with uniform delay step ( $d\tau$ ). Convert the energy axis to frequency by  $f = \epsilon/2\pi$ . The frequency start and end points are  $f_1$  and  $f_N$ , respectively. The central frequency is:

$$f_0 = \frac{f_N - f_1}{2} \quad (3)$$

**Step 2:** Before fed to reconstruction algorithm, the measured spectrogram should be modified to:

$$I_a(p, \tau) = \frac{I_0(p, \tau)}{|d(p)|^2} \quad (4)$$

**Why to do it?**

With this modification, the spectrogram becomes:

$$I_a(p, \tau) = \left| \int_{-\infty}^{\infty} E_X(t - \tau) \frac{d(p + A(t))}{d(p)} e^{-i\varphi(p,t)} e^{\left(\frac{p^2}{2} + I_p\right)t} dt \right|^2 \quad (5)$$

By applying *central momentum approximation*,  $p$  is substituted with  $p_0$  in  $d$  and  $\varphi$ , so

$$I_a(p, \tau) \approx \left| \int_{-\infty}^{\infty} E_X(t - \tau) \frac{d(p_0 + A(t))}{d(p_0)} e^{-i\varphi(p_0,t)} e^{\left(\frac{p^2}{2} + I_p\right)t} dt \right|^2 \quad (6)$$

$$\triangleq \left| \int_{-\infty}^{\infty} E_X(t - \tau) G(t) e^{\left(\frac{p^2}{2} + W\right)t} dt \right|^2 \quad (7)$$

here,

$$G(t) = \frac{d(p_0 + A(t))}{d(p_0)} e^{-i\varphi(p_0,t)} \quad (8)$$

**Step 3:** Interpolating  $I_a$  along frequency axis to change the number of frequency points from  $N_\epsilon$  to  $M_\epsilon$  so that:

1.  $M_\epsilon$  is a power of two;
2. the frequency points are uniformly spaced with interval of  $df$ :

$$df = \frac{f_N - f_1}{M_\epsilon}; \quad (9)$$

3.  $df$  satisfies:

$$df d\tau = L/M_\epsilon, \quad (10)$$

here  $L$  is an integer.

**Why to do it?**

After Interpolation, the IAP pulse  $P(t)$  is then digitized at equally spaced time points with time step of  $dt$ :

$$dt = \frac{1}{df M_\epsilon} \quad (11)$$

Eq. (11) ensures that the energy points of  $I_a$  matches the Fourier transform of  $P(t)$ . Combining Eq. (10) and Eq. (11), we have:

$$d\tau = L dt, \quad (12)$$

so Eq. (10) requires that the delay interval should be integer multiples of  $dt$ .

**How to do it?**

There may be no solutions for  $M_\varepsilon$  to fulfill the three conditions simultaneously. However, the interpolation can be done with the following procedure:

1. set  $L$  to a proper integer so that:

$$L \geq d\tau(f_N - f_1) \quad (13)$$

2.  $dt$  can then be calculated as:

$$dt = \frac{d\tau}{L}; \quad (14)$$

3. set  $M_\varepsilon$  to be a power of two. Note that  $M_\varepsilon$  cannot be too small since it determines the time window of the IAP pulse by  $T = (M_\varepsilon - 1)dt$ :

4. the frequency interval is:

$$df = \frac{1}{M_\varepsilon dt} \quad (15)$$

5. generate the sampling frequency points by keep the central frequency unchanged:

$$f_1 = f_0 - \frac{M_\varepsilon}{2} df, \quad (16)$$

$$f_N = f_0 + \left(\frac{M_\varepsilon}{2} - 1\right) df, \quad (17)$$

$$(18)$$

so,

$$f_i = f_0 + \left(i - 1 - \frac{M_\varepsilon}{2}\right) df \quad (i = 1, 2, \dots, M_\varepsilon) \quad (19)$$

- 6.